Spin scale-discretised wavelets on the sphere for the analysis of CMB polarisation

Jason McEwen www.jasonmcewen.org @jasonmcewen

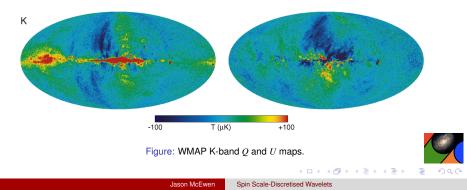
Mullard Space Science Laboratory (MSSL) University College London (UCL)

In collaboration with Martin Büttner, Boris Leistedt, Hiranya Peiris, Pierre Vandergheynst & Yves Wiaux

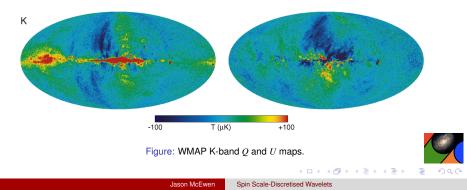
7th International Conference of the ERCIM on Computational and Methodological Statistics Pisa, December 2014

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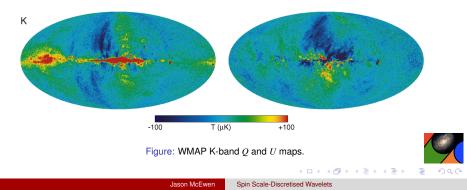
- CMB polarisation contains a wealth of cosmological information; but extracting this information is challenging.
- Evidence for primordial gravitational waves from BICEP2?
- Observe *Q* and *U* Stokes parameters.
- Construct $Q \pm iU$, which is a spin ± 2 field: $(Q \pm iU)'(\omega) = e^{\pm i2\chi}(Q \pm iU)(\omega)$.
- Extract and analyse cosmological maps.



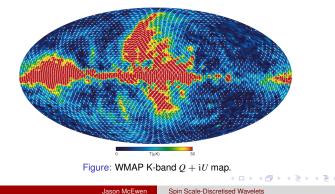
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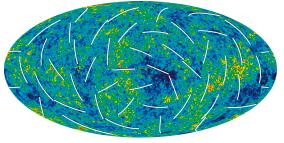


Figure: WMAP cosmological temperature and polarisation.



Outline

Spin scale-discretised wavelets on the sphere

- Past algorithms
- E/B separation



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Recall wavelet transform in Euclidean space

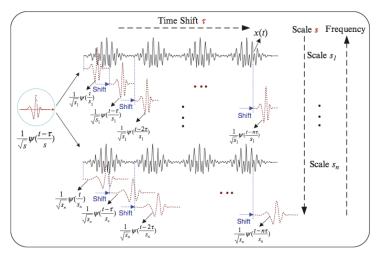


Figure: Wavelet scaling and shifting [Credit: Gao & Yan (2010)]



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Wavelets on the sphere

Dilation and translation

- Construct wavelet atoms from affine transformations (dilation, translation) on the sphere of a mother wavelet.
- The natural extension of translations to the sphere are rotations. Rotation of a function *f* on the sphere is defined by

$$[\mathcal{R}(\rho)f](\omega) = f(\mathsf{R}_{\rho}^{-1} \cdot \omega), \quad \omega = (\theta, \varphi) \in \mathbb{S}^{2}, \quad \rho = (\alpha, \beta, \gamma) \in \mathrm{SO}(3) \; .$$

translation

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• How define dilation on the sphere?

- Stereographic projection Antoine & Vandergheynst (1999), Wiaux et al. (2005)
- Harmonic dilation wavelets McEwen *et al.* (2006), Sanz *et al.* (2006)
- Isotropic undecimated wavelets Starck *et al.* (2005), Starck *et al.* (2009)
- Needlets Narcowich et al. (2006), Baldi et al. (2009), Marinucci et al. (2008), Geller et al. (2008)
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Wavelet construction

- Exact reconstruction with directional wavelets on the sphere Wiaux, McEwen, Vandergheynst, Blanc (2008)
- Extend to spin functions.

 Scale-discretised wavelet sΨ^j ∈ L²(S²) defined in harmonic space:

$${}_{s}\Psi^{j}_{\ell m} \equiv \kappa^{j}(\ell) s_{\ell m} \, .$$

• Admissible wavelets constructed to satisfy a resolution of the identity:

$$\begin{split} \boxed{|_{s}\Phi_{\ell 0}|^{2}} + \sum_{j=0}^{J} \sum_{m=-\ell}^{\ell} \boxed{|_{s}\Psi_{\ell m}^{j}|^{2}}_{\text{wavelet}} = 1, \quad \forall \ell. \end{split}$$



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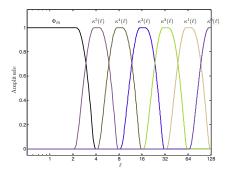


Figure: Harmonic tiling on the sphere.

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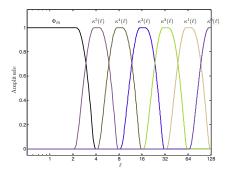


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$$\begin{split} & \underbrace{|_{s} \Phi_{\ell 0}|^{2}}_{\text{scaling function}} + \sum_{j=0}^{J} \sum_{m=-\ell}^{\ell} \underbrace{|_{s} \Psi_{\ell m}^{j}|^{2}}_{\text{wavelet}} = 1, \quad \forall \ell. \end{split}$$

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Spin Wavelets Fast Algorithms E/B Separation

Scale-discretised wavelets on the sphere Wavelets

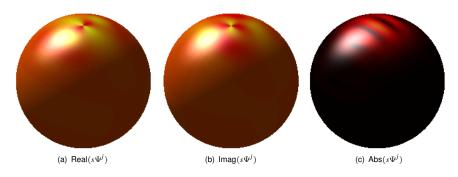


Figure: Spin scale-discretised wavelets on the sphere.



Forward and inverse transform (i.e. analysis and synthesis)

• The spin scale-discretised wavelet transform is given by the usual projection onto each wavelet:

$$\frac{W^{s\Psi^{j}}(\rho) = \langle_{s}f, \mathcal{R}_{\rho s}\Psi^{j}\rangle}{\text{projection}} = \int_{\mathbb{S}^{2}} d\Omega(\omega)_{s}f(\omega)(\mathcal{R}_{\rho s}\Psi^{j})^{*}(\omega) .$$

- Wavelet coefficients are scalar and not spin.
- Wavelet coefficients live in $SO(3) \times \mathbb{Z}$; thus, directional structure is naturally incorporated.
- The original function may be recovered exactly in practice from the wavelet (and scaling) coefficients:

$$sf(\omega) = \sum_{j=0}^{J} \int_{SO(3)} d\varrho(\rho) W^{s\Psi^{j}}(\rho) (\mathcal{R}_{\rho \ s} \Psi^{j})(\omega) .$$
finite sum wavelet contribution



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Fast Wigner transform

Fourier transform on the rotation group SO(3)

- Wavelet coefficients live on the rotation group SO(3): $W^{s\Psi^{j}} \in L^{2}(SO(3))$.
- Develop fast wavelet transforms by considering their (Wigner) harmonic representation.
- Signal on the rotation group $F \in L^2(SO(3))$ may expressed by Wigner decomposition:

$$F(\rho) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \sum_{n=-\ell}^{\ell} \frac{2\ell+1}{8\pi^2} F_{nm}^{\ell} D_{nm}^{\ell*}(\rho)$$

where Wigner coefficients given by usual projection onto basis functions:

$$F_{mn}^{\ell} = \langle F, D_{mn}^{\ell*} \rangle = \int_{\mathrm{SO}(3)} \mathrm{d}\varrho(\rho) F(\rho) D_{mn}^{\ell}(\rho) \,.$$

Novel sampling theorem on the rotation group requiring $2L^3$ samples only

to represent a signal band-limited at L (follows straightforwardly from McEwen & Wiaux 2012)



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Fast spin scale-discretised wavelet transform

Exact and efficient computation via Wigner transforms

• Wavelet analysis can be posed as an inverse Wigner transform on SO(3):

$$\left(W^{\Psi^{j}}\right)_{mn}^{\ell} = \frac{8\pi^{2}}{2\ell+1} f_{\ell m} \Psi_{\ell n}^{j*} .$$
analysis

• Wavelet synthesis can be posed as a forward Wigner transform on SO(3):

$$f(\omega) = \sum_{j=0}^{J} \sum_{\ell m n} \frac{2\ell+1}{8\pi^2} \left(W^{\Psi^j} \right)_{m n}^{\ell} \Psi^j_{\ell n} Y_{\ell m}(\omega) ,$$

synthesis

where

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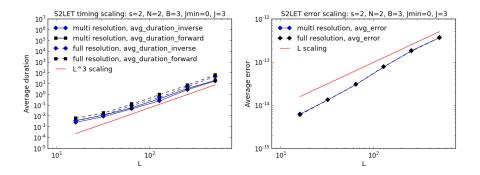
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Spin Wavelets Fast Algorithms E/B Separation

Fast spin scale-discretised wavelet transform Timing and accuracy





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E/B separation From $Q \pm iU$ to E and B maps

- Different physical processes exhibit different symmetries and thus behave differently under parity transformation.
- Can exploit this property to separate signals arising from different underlying physical mechanisms.
- Decompose $Q \pm iU$ into parity even and odd components:

$$\tilde{E}(\omega) = -\frac{1}{2} \Big[\bar{\eth}^2 (Q + \mathrm{i}U)(\omega) + \eth^2 (Q - \mathrm{i}U)(\omega) \Big] \Big\|_{\mathrm{H}}^{\mathrm{pp}}$$

$$\tilde{B}(\omega) = \frac{\mathrm{i}}{2} \Big[\tilde{\eth}^2 (\mathcal{Q} + \mathrm{i}U)(\omega) - \eth^2 (\mathcal{Q} - \mathrm{i}U)(\omega) \Big] \Big] \begin{bmatrix} \mathbf{q} \\ \mathbf{q} \\ \mathbf{q} \end{bmatrix}$$

where $\bar{\eth}$ and \eth are spin lowering and raising operators, respectively.

• Number of existing techniques: Lewis *et al.* (2002), Bunn *et al.* (2003), Bowyer *et al.* (2011), Kim (2013).



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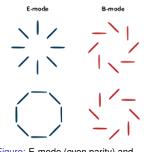


Figure: E-mode (even parity) and B-mode (odd parity) signals [Credit:

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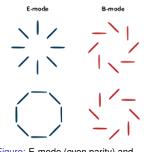


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E/B separation

Spin and scalar scale-discretised wavelets

• Spin wavelet transform of Q + iU (observable):

$$\frac{W_{Q+iU}^{2\Psi^{j}}(\rho) = \langle Q + iU, \mathcal{R}_{\rho} \,_{2}\Psi^{j} \rangle}{\text{spin wavelet transform}} = \int_{\mathbb{S}^{2}} \mathrm{d}\Omega(\omega)(Q + iU)(\omega)(\mathcal{R}_{\rho} \,_{2}\Psi^{j})^{*}(\omega) \,.$$

• Scalar wavelet transforms of E and B (non-observable):

$$\begin{split} W_{\tilde{E}}^{0\tilde{\Psi}^{j}}(\rho) &= \langle \tilde{E}, \ \mathcal{R}_{\rho \ 0}\tilde{\Psi}^{j} \rangle \\ \text{scalar wavelet transform} \\ \hline W_{\tilde{B}}^{0\tilde{\Psi}^{j}}(\rho) &= \langle \tilde{B}, \ \mathcal{R}_{\rho \ 0}\tilde{\Psi}^{j} \rangle \\ \text{scalar wavelet transform} \\ = \int_{\mathbb{S}^{2}} \mathrm{d}\Omega(\omega)\tilde{E}(\omega)(\mathcal{R}_{\rho \ 0}\tilde{\Psi}^{j})^{*}(\omega) , \end{split}$$

where $_{0}\tilde{\Psi}^{j}\equiv\bar{\eth}^{2}{}_{2}\Psi^{j}.$

Spin wavelet coefficients of Q + iU are connected to scalar wavelet coefficients of E/B:

$$W_{\tilde{E}}^{\tilde{\Psi}^{j}}(\rho) = -\mathrm{Re}\Big[W_{Q+\mathrm{i}U}^{2\Psi^{j}}(\rho)\Big] \quad \mathrm{and} \quad W_{\tilde{B}}^{0\tilde{\Psi}^{j}}(\rho) = -\mathrm{Im}\Big[W_{Q+\mathrm{i}U}^{2\Psi^{j}}(\rho)\Big] \; .$$



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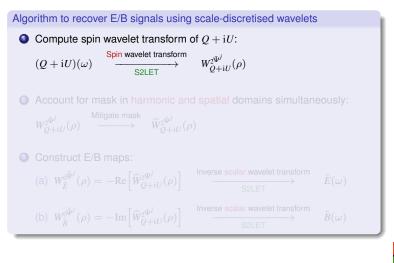
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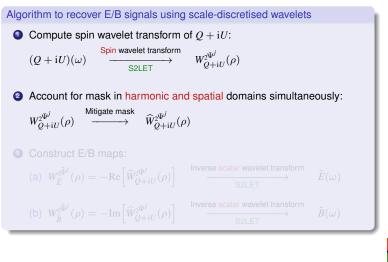
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E/B separation Using scale-discretised wavelets





E/B separation Using scale-discretised wavelets



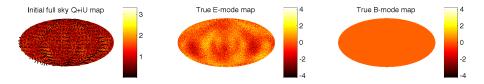


E/B separation Using scale-discretised wavelets

Algorithm to recover E/B signals using scale-discretised wavelets • Compute spin wavelet transform of O + iU: Spin wavelet transform $W^{2\Psi^j}_{O+\mathrm{i}U}(\rho)$ $(O + iU)(\omega)$ S2LET Account for mask in harmonic and spatial domains simultaneously: $W^{2\Psi^{j}}_{O+iU}(\rho) \xrightarrow{\text{Mitigate mask}} \widehat{W}^{2\Psi^{j}}_{O+iU}(\rho)$ Onstruct E/B maps: Inverse scalar wavelet transform (a) $W_{\tilde{E}}^{0\tilde{\Psi}^{j}}(\rho) = -\operatorname{Re}\left[\widehat{W}_{O+\mathrm{i}U}^{2\Psi^{j}}(\rho)\right]$ $\tilde{E}(\omega)$ S2LET Inverse scalar wavelet transform (b) $W^{0\tilde{\Psi}^{j}}_{\tilde{B}}(\rho) = -\mathrm{Im}\left[\widehat{W}^{2\Psi^{j}}_{Q+\mathrm{i}U}(\rho)\right]$ $\tilde{B}(\omega)$ S2LET



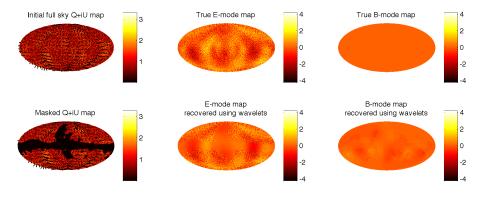
E/B separation Simulations





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E/B separation Simulations





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0.15

0.1

0.05

0.3 0.25 0.2 0.15 0.1 0.05

E/B separation Preliminary results

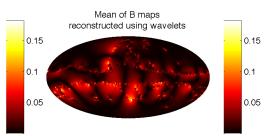
Mean of B maps reconstructed using harmonics Std dev of B maps reconstructed using harmonics

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E/B separation Preliminary results

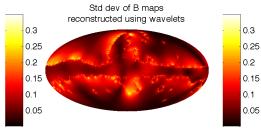
Mean of B maps reconstructed using harmonics





Std dev of B maps reconstructed using harmonics





Summary

Spin scale-discretised wavelets are a powerful tool to study CMB polarisation, with an elegant and practical connection to scalar wavelet transforms of E/B maps.



