Sparsity, Euclid and the SKA

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Big cosmology: big science, big data and big algos

- A new era of big cosmology is emerging.
 - Planck: full-sky observations of the CMB at unprecedented resolution, sensitivity and frequency coverage.
 - Euclid: unprecedented survey of billion galaxies over more than one third of the sky.
 - Square Kilometre Array (SKA): sensitivity 50x that of previous radio telescopes with phenomenal data rates.
 - Others...



(a) Planck

(b) Euclid

(c) SKA

 New instruments must be complemented with novel analyses methodologies to extract new science from big data-sets

$$\rightarrow$$
 sparsity.

Outline







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Outline



2 Sparsity and Euclid

Sparsity and the SKA

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What is sparsity?

What is sparsity?

- representation of data in such a way that many data points are zero.

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What is sparsity?



What is sparsity?



Sparsifying transform





Why is sparsity useful?

Why is sparsity useful?

- efficient characterisation of information.

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Why is sparsity useful?



Add noise





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Why is sparsity useful?



Sparsifying transform



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Why is sparsity useful?



Sparsifying transform





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Why is sparsity useful?



Threshold





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Why is sparsity useful?



Inverse transform





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Why is sparsity useful?



(a) Original



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[Credit: http://www.ceremade.dauphine.fr/~peyre/numerical-tour/tours/denoisingwav_2_wavelet_2d/]

⁽b) Noisy

Outline







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How can we construct sparsifying transforms?

How can we construct sparsifying transforms?

- many signals in nature have spatially localised, scale-dependent features.

How can we construct sparsifying transforms?







Haar (1909) Morlet and Grossman (1981)



Figure: Fourier vs wavelet transform [Credit: http://www.wavelet.org/tutorial/]

How can we construct sparsifying transforms?







Haar (1909) Morlet and Grossman (1981)



Figure: Fourier vs wavelet transform [Credit: http://www.wavelet.org/tutorial/]

How can we construct sparsifying transforms?



Cosmological observations made on celestial sphere

- Cosmological observations are inherently made on the celestial sphere.
 - Observations of the cosmic microwave background (CMB) are made on the sphere.
 - Observations tracing the large-scale structure (LSS) are made on the ball.





(a) CMB (WMAP)



Scale-discretised wavelets on the sphere

- Exact reconstruction with directional wavelets on the sphere Wiaux, McEwen, Vandergheynst, Blanc (2008) [arXiv:0712.3519]
- Alternatives: isotropic wavelets, pyramidal wavelets, ridgelets, curvelets (Starck et al. 2006); needlets (Narcowich et al. 2006, Baldi et al. 2009, Marinucci et al. 2008).



Figure: Harmonic tiling on the sphere. $\langle \Box \rangle$ $\langle \Box \rangle$

Scale-discretised wavelets on the sphere



Figure: Scale-discretised wavelets on the sphere.

• The scale-discretised wavelet transform is given by the usual projection onto each wavelet:

$$W^{\Psi^j}(\rho) \equiv (f \star \Psi^j)(\rho) = \langle f, \ \mathcal{R}_\rho \Psi^j \rangle = \int_{\mathbb{S}^2} \mathrm{d}\Omega(\omega) f(\omega) (\mathcal{R}_\rho \Psi^j)^*(\omega) \ ,$$

 The original function may be recovered exactly in practice from the wavelet (and scaling) coefficients:

$$f(\omega) = 2\pi \int_{\mathbb{S}^2} \mathrm{d}\Omega(\omega') W^{\Phi}(\omega')(\mathcal{R}_{\omega'} L^{\mathrm{d}} \Phi)(\omega) + \sum_{j=0}^J \int_{\mathrm{SO}(3)} \mathrm{d}\varrho(\rho) W^{\Psi^j}(\rho)(\mathcal{R}_{\rho} L^{\mathrm{d}} \Psi^j)(\omega) \,.$$

Codes for scale-discretised wavelets on the sphere



S2DW code

http://www.s2dw.org

Exact reconstruction with directional wavelets on the sphere Wiaux, McEwen, Vandergheynst, Blanc (2008) [arXiv:0712.3519]

- Fortran
- Parallelised
- Supports directional, steerable wavelets

S2LET code

http://www.s2let.org

S2LET: A code to perform fast wavelet analysis on the sphere Leistedt, McEwen, Vandergheynst, Wiaux (2012) [arXiv:1211.1680]

- C, Matlab, IDL, Java
- Supports only axisymmetric wavelets at present
- Future extensions:
 - Directional, steerable wavelets
 - Faster algorithms to perform wavelet transforms
 - Spin wavelets



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Fourier-LAGuerre wavelets (flaglets) on the ball



Figure: Tiling of Fourier-Laguerre space.

 Exact wavelets on the ball Leistedt & McEwen (2012) [arXiv:1205.0792]

- Extend scale-discretised wavelets on the sphere to the ball.
- Some subtleties (define translation and convolution on the radial line).
- Construct wavelets by tiling the ℓ -*p* harmonic plane.

• The Fourier-Laguerre wavelet transform is given by the usual projection onto each wavelet:

$$W^{\Psi^{jj'}}(\mathbf{r}) \equiv (f \star \Psi^{jj'})(\mathbf{r}) = \langle f | \mathcal{T}_{\mathbf{r}} \Psi^{jj'} \rangle_{\mathbb{B}^3} = \int_{B^3} \mathrm{d}^3 \mathbf{r}' f(\mathbf{r}') (\mathcal{T}_{\mathbf{r}} \Psi^{jj'})(\mathbf{r}') \; .$$

• The original function may be synthesised exactly in practice from its wavelet (and scaling) coefficients:

$$f(\mathbf{r}) = \int_{B^3} \mathrm{d}^3 \mathbf{r}' W^{\Phi}(\mathbf{r}') (\mathcal{T}_{\mathbf{r}} \Phi)(\mathbf{r}') + \sum_{j=J_0}^J \sum_{j'=J_0'}^{J'} \int_{B^3} \mathrm{d}^3 \mathbf{r}' W^{\Psi_{j}^{jj'}}(\mathbf{r}') (\mathcal{T}_{\mathbf{r}} \Psi^{jj'})(\mathbf{r}') \; .$$

Alternatives: Spherical 3D isotropic wavelets (Lanusse, Rassat & Starck 2012)

Fourier-LAGuerre wavelets (flaglets) on the ball



Figure: Scale-discretised wavelets on the ball.

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Codes for Fourier-LAGuerre wavelets (flaglets) on the ball



FLAG code: Fourier-Laguerre transform

http://www.flaglets.org

Exact wavelets on the ball Leistedt & McEwen (2012) [arXiv:1205.0792]

- C, Matlab, IDL, Java
- Exact Fourier-LAGuerre transform on the ball



FLAGLET code: Fourier-Laguerre wavelets

http://www.flaglets.org

Exact wavelets on the ball Leistedt & McEwen (2012) [arXiv:1205.0792]

- C, Matlab, IDL, Java
- Exact (Fourier-LAGuerre) wavelets on the ball coined flaglets!

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Large-scale structure (LSS) of the Universe

• Map Horizon simulation of large-scale structure (LSS) to Fourier-Laguerre sampling.





LSS fly through

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Flaglet void finding

- Find voids in the large-scale structure (LSS) of the Universe.
- Perform Alcock & Paczynski (1979) test: study void shapes to constrain the nature of dark energy (e.g. Sutter et al. 2012).

LSS voids

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Outline



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- Next evolution of wavelet analysis \rightarrow wavelets are a key ingredient.
- The mystery of JPEG compression (discrete cosine transform; wavelet transform).
- Move compression to the acquisition stage \rightarrow compressive sensing.
- Deep mathematical foundation (Candes et al. 2006, Donoho 2006).



Figure: Single pixel camera

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• Linear operator (algebra) representation of signal decomposition (into atoms of a dictionary):

$$x(t) = \sum_{i} \alpha_{i} \Psi_{i}(t) \quad \rightarrow \quad \mathbf{x} = \sum_{i} \Psi_{i} \alpha_{i} = \begin{pmatrix} | \\ \Psi_{0} \\ | \end{pmatrix} \alpha_{0} + \begin{pmatrix} | \\ \Psi_{1} \\ | \end{pmatrix} \alpha_{1} + \cdots \quad \rightarrow \quad \boxed{\mathbf{x} = \Psi \boldsymbol{\alpha}}$$

• Linear operator (algebra) representation of measurement:

$$y_i = \langle x, \Phi_j \rangle \rightarrow \mathbf{y} = \begin{pmatrix} -\Phi_0 & -\\ -\Phi_1 & -\\ \vdots \end{pmatrix} \mathbf{x} \rightarrow \mathbf{y} = \Phi \mathbf{x}$$

• Putting it together:

$$\mathbf{y} = \Phi \mathbf{x} = \Phi \Psi \boldsymbol{\alpha}$$



• Ill-posed inverse problem:

$$y = \Phi x + n = \Phi \Psi \alpha + n.$$

 Solve by imposing a regularising prior that the signal to be recovered is sparse in Ψ, *i.e.* solve the following ℓ₀ optimisation problem:

$$\boldsymbol{\alpha}^{\star} = \operatorname*{arg\,min}_{\boldsymbol{\alpha}} \| \boldsymbol{\alpha} \|_{0} \ \, \text{such that} \ \, \| \boldsymbol{y} - \Phi \Psi \boldsymbol{\alpha} \|_{2} \leq \epsilon \ ,$$

where the signal is synthesising by $x^{\star} = \Psi \alpha^{\star}$.

• Recall norms given by:

$$\|\alpha\|_0 =$$
no. non-zero elements $\|\alpha\|_1 = \sum_i |\alpha_i| \qquad \|\alpha\|_2 = \left(\sum_i |\alpha_i|^2\right)^{1/2}$

• Solving this problem is difficult (combinatorial).

• Instead, solve the ℓ_1 optimisation problem (convex):

 $oldsymbol{lpha}^{\star} = \operatorname*{arg\,min}_{oldsymbol{lpha}} \|lpha\|_1 \, \, ext{such that} \, \, \|\mathbf{y} - \Phi \Psi oldsymbol{lpha}\|_2 \leq \epsilon \, .$

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$$\|\alpha\|_0 = \text{no. non-zero elements}$$
 $\|\alpha\|_1 = \sum_i |\alpha_i|$ $\|\alpha\|_2 = \left(\sum_i |\alpha_i|^2\right)^{1/2}$

- Solving this problem is difficult (combinatorial).
- Instead, solve the ℓ_1 optimisation problem (convex):

$$\boldsymbol{\alpha}^{\star} = \underset{\boldsymbol{\alpha}}{\arg\min} \|\boldsymbol{\alpha}\|_{1} \text{ such that } \|\boldsymbol{y} - \boldsymbol{\Phi} \boldsymbol{\Psi} \boldsymbol{\alpha}\|_{2} \leq \epsilon \ .$$

- The solutions of the ℓ_0 and ℓ_1 problems are often the same.
- Geometry of ℓ_2 and ℓ_1 problems.



Figure: Geometry of (a) ℓ_0 (b) ℓ_2 and (c) ℓ_1 problems. [Credit: Baraniuk (2007)]

Compressive sensing

- In the absence of noise, compressed sensing is exact!
- Number of measurements required to achieve exact reconstruction is given by

$M \ge c\mu^2 K \log N ,$

where *K* is the sparsity and *N* the dimensionality.

• The coherence between the measurement and sparsity basis is given by



$\mu = \sqrt{N} \max_{i,j} |\langle \Psi_i, \Phi_j \rangle|$.

Robust to noise.

Radio interferometric inverse problem

• Consider the ill-posed inverse problem of radio interferometric imaging:

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y = \Phi x + n \, ,
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where *y* are the measured visibilities, Φ_p is the linear measurement operator, x_p is the underlying image and *n* is instrumental noise.

- Measurement operator $\Phi = \mathbf{M} \mathbf{F} \mathbf{C} \mathbf{A}$ may incorporate:
 - primary beam A of the telescope;
 - w-component modulation C (responsible for the spread spectrum phenomenon);
 - Fourier transform F;
 - masking M which encodes the incomplete measurements taken by the interferometer.

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Radio interferometric imaging

 SARA algorithm for radio interferometric imaging, building on compressive sensing techniques (Carrillo, McEwen & Wiaux 2012) [arXiv:1205.3123].



Continuous visibilities

• PURIFY: realistic radio interferometric imaging with compressive sensing (Carrillo, McEwen & Wiaux 2013) [arXiv:1307.4370]. http://basp-group.github.io/purify/





Jason McEwen

Sparsity, Euclid and the SKA

Wide field-of-view

- Wide fields give rise to the spread spectrum effect (Wiaux *et al.* 2009), which improves reconstruction quality.
- Recently studied in a more realistic setting (Wolz, McEwen, Abdalla, Carrillo, Wiaux 2013) [arXiv:1307.3424].



(a) No spread spectrum



(C) Idealised spread spectrum

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Figure: Reconstruction fidelity in the presence and absence of the spread spectrum effect.

Wide field-of-view

- Wide fields give rise to the spread spectrum effect (Wiaux *et al.* 2009), which improves reconstruction quality.
- Recently studied in a more realistic setting (Wolz, McEwen, Abdalla, Carrillo, Wiaux 2013) [arXiv:1307.3424].



(a) No spread spectrum (b) More realistic spread spectrum

(C) Idealised spread spectrum

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Figure: Reconstruction fidelity in the presence and absence of the spread spectrum effect.

Summary

For big cosmology we need novel analysis methods to deal with the data deluge of forthcoming experiments (*e.g.* Euclid, SKA, ...)

 \rightarrow exploit sparsity.